

LES of the Ocean Mixed Layer with its Application to the Parameterization in a Climate Model

Yign Noh

*Department of Atmospheric Sciences/Global Environmental Laboratory,
Yonsei University, Korea*

Parameterization of SGS Turbulence in the Climate Model

- Grid sizes of the climate model (*even at Earth Simulator*) are still much larger than the eddy size of turbulence
 - ~ 10-100 km (horizontal)
 - ~ 10-100 m (vertical)

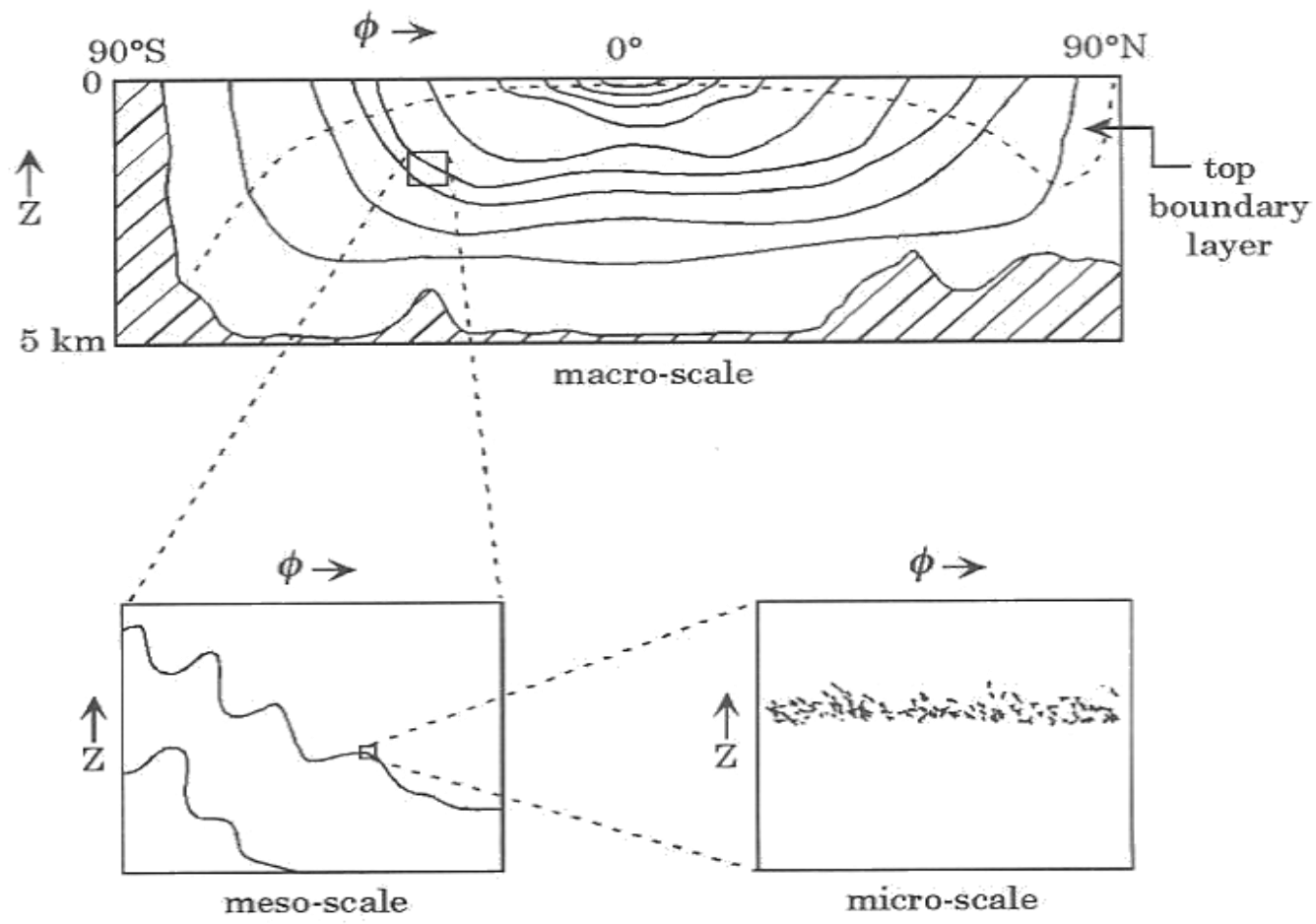
⇒ *We have to parameterize the turbulence within the grid.*

ex) OGCM - ocean mixed layer, ocean deep convection,
meso-scale eddies,

AGCM – PBL , cumulus convection, cloud

Biogeochemical Model

- particle flux, aerosol, biogeochemical process



How is the SGS parameterization of turbulence developed currently ?

- A new parameterization based on **intuitive** (or *arbitrary*) hypotheses
 - Show that the new parameterization performs better **in a certain case**
 - comparison with observation data
 - examination of the performance of a climate model
 - **Claim** to improve the parameterization (and *publish*).

- There is no guarantee that
 - *We are really moving to the right direction for the realistic representation of nature ultimately.*

(Driving without a map)

 - *The hypotheses really represent correct physics.*

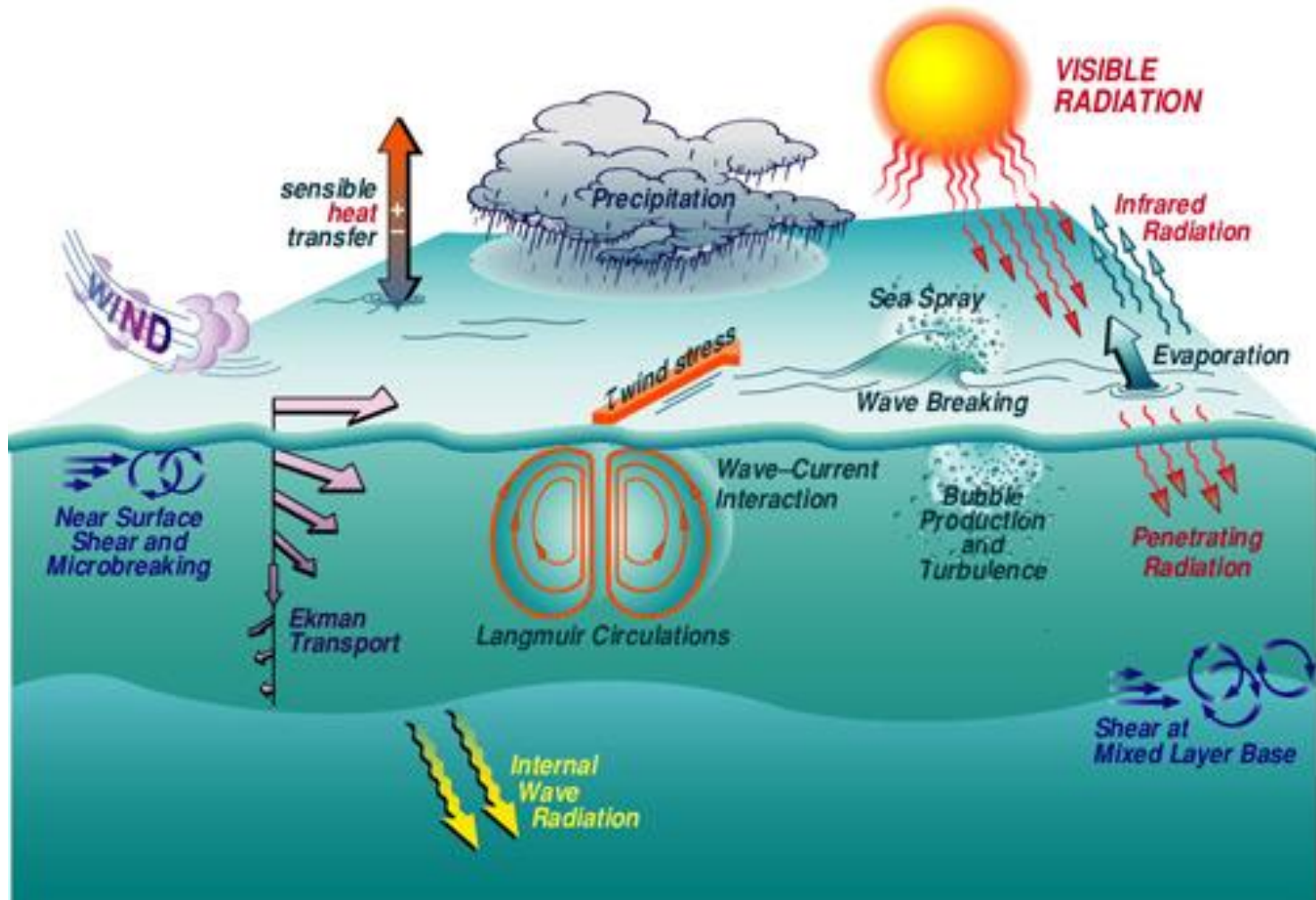
LES (large eddy simulation), which reproduces realistic 3-D structure of turbulence,

provides a powerful tool for the development of the SGS parameterization of turbulence

by examining

- not only *the performance of the parameterization*
- but also *the hypotheses* used in the parameterization.

Ocean Mixed Layer



LES of Ocean Mixed Layer

- Parallel Large Eddy Simulation Model (PALM)
- Inclusion of the effects of *Langmuir circulation* and *wave breaking*
(Noh et al. , *JPO* 2004)

$$\frac{\partial u_i}{\partial t} + (u_j + u_{sj}) \frac{\partial u_i}{\partial x_j} = -\frac{1}{\rho_0} \frac{\partial p}{\partial x_i} + \nu \frac{\partial^2 u_i}{\partial x_j \partial x_j}$$

$$- \varepsilon_{ijk} f_j u_k + u_{sk}$$

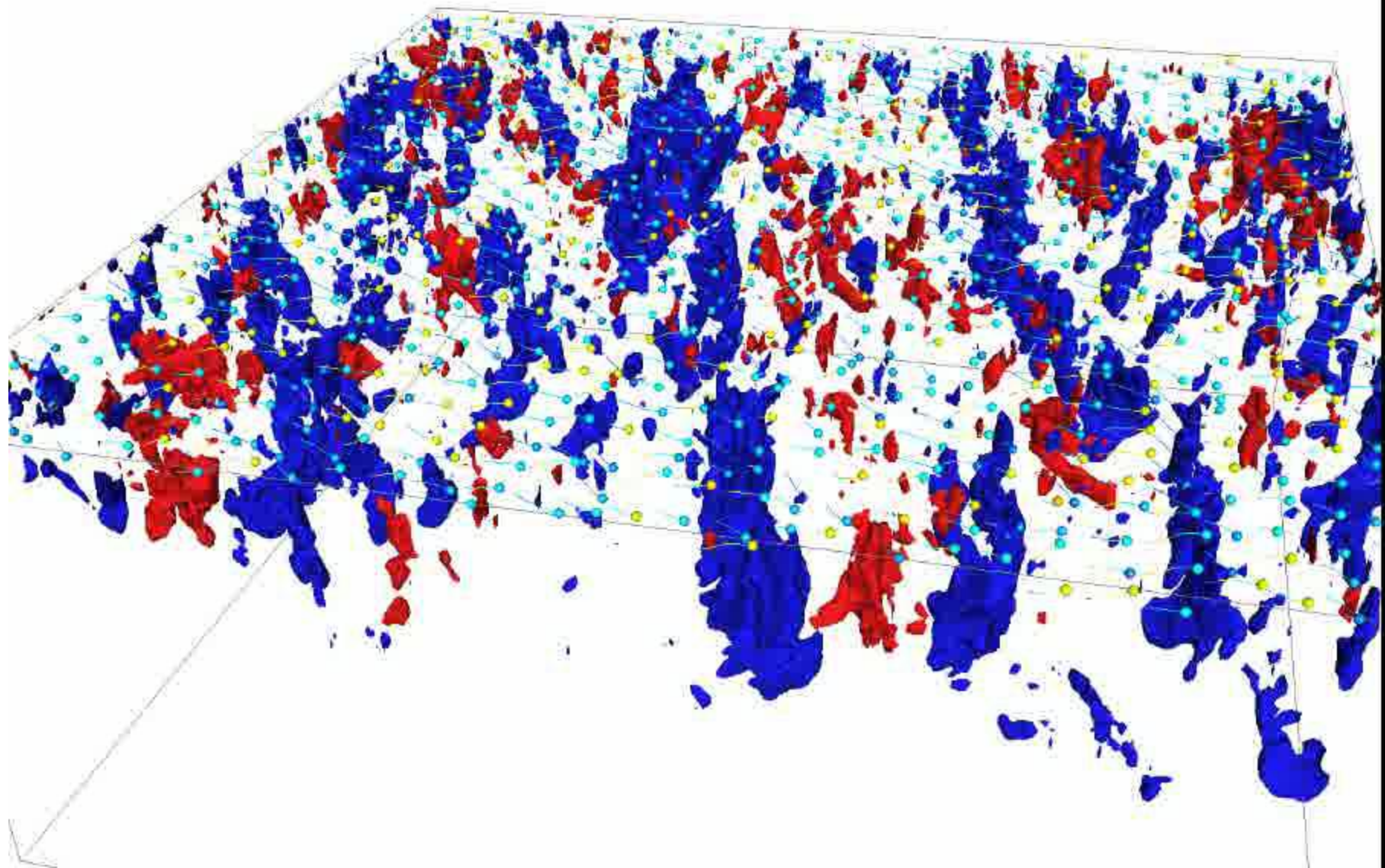
$$+ \varepsilon_{ijk} u_{sj} \omega_k \quad \leftarrow \text{vortex force to generate LC}$$

$$+ F_i(\mathbf{x}, t) \quad \leftarrow \text{stochastic forcing to generate wave breaking}$$

$$* \quad F_i = \frac{\alpha u_*}{\tau_0} G(x, y, t) \delta(z) \delta_{i3} - \delta_{i3}$$

$$u_s = U_s \exp(-4\pi z / \lambda)$$

$$U_s = (2\pi a / \lambda)^2 (g\lambda / 2\pi)^{1/2}$$



1. What is the Role of **Langmuir Circulation** in the Mixed Layer Deepening ?
2. What is the Role of **Wind Stress** in the Convective Deepening of the Mixed Layer

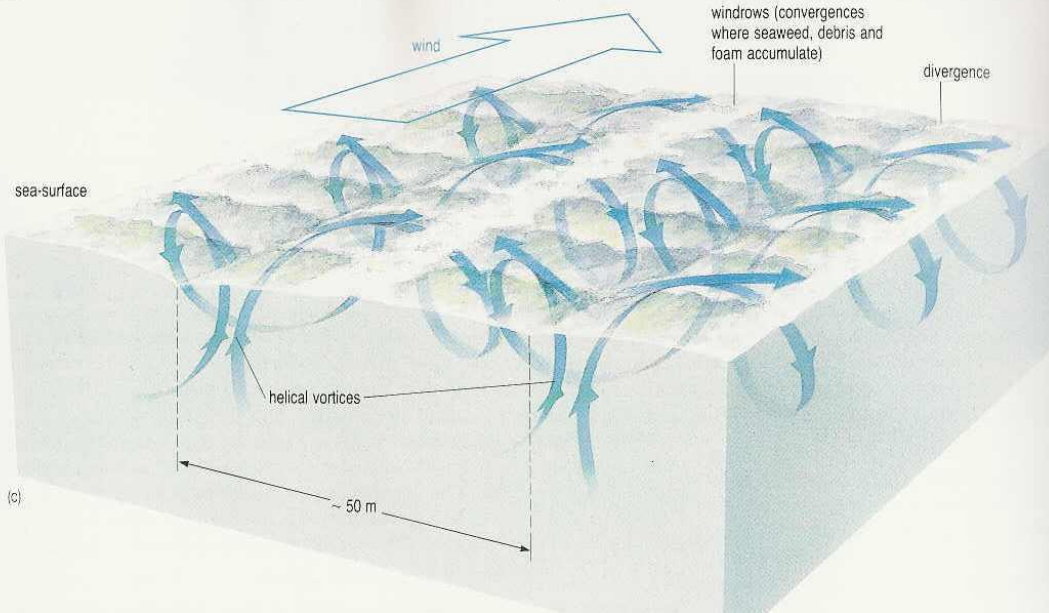
1. What is the Role of **Langmuir Circulation** in the Mixed Layer Deepening ?



(a)



(b)



(c)

Langmuir circulation observed in the ocean

Influence of LC on Vertical Mixing

- It has been well known that LC enhances vertical mixing greatly within the mixed layer.
- The role in the mixed layer deepening is still under debate.
 - * Li et al. (1995), Sullivan et al. (2007), Grant and Belcher (2009), Kukulka et al. (2009)
 - enhanced mixed layer deepening
 - * Weller & Price (1988), Thorpe et al. (2003)
 - no evidence of the contribution by LC
 - * Skillingstad et al. (2000)
 - The effects of LC are mostly confined to the initial stage of mixed layer growth.

Consideration of LC in the Mixed Layer Model

- Li et al. (1995), Smith (1998)
 - modification of the PWP model

*Mixing occurs, if

$$h\Delta B / (\Delta U)^2 < 0.65 \quad \leftarrow \text{PWP model}$$

$$h\Delta B / v_L^2 < C \quad (C = 9.8)$$

$$v_L = (U_s u_*^2)^{1/3} \quad : \text{velocity scale of LC}$$

- McWilliams & Sullivan (2000), Smyth et al. (2002)
 - modification of the KPP model

$$w = \frac{\kappa u_*}{\phi(z,t)} \quad \rightarrow \quad w = \frac{\kappa u_*}{\phi(z,t)} \left[1 + \frac{C_w(u_*, w_*)}{La^4} \right]^{1/2}$$

$$-\overline{uw} = K_m \frac{\partial U}{\partial z} \quad \rightarrow \quad -\overline{uw} = K_m \left(\frac{\partial U}{\partial z} + \gamma_u \right)$$

$$K(z,t) = h(t)w(z,t)G(z,t)$$

- D'Alescio et al. (1998)
 - modification of TKE equation and BC at the MLD

$$\frac{\partial E}{\partial t} = F + P_s - P_b - \varepsilon + P_L$$

$$P_L = u_s \partial(\overline{u'w'}) / \partial z \quad : \text{additional TKE production by the vortex force}$$

Experiments

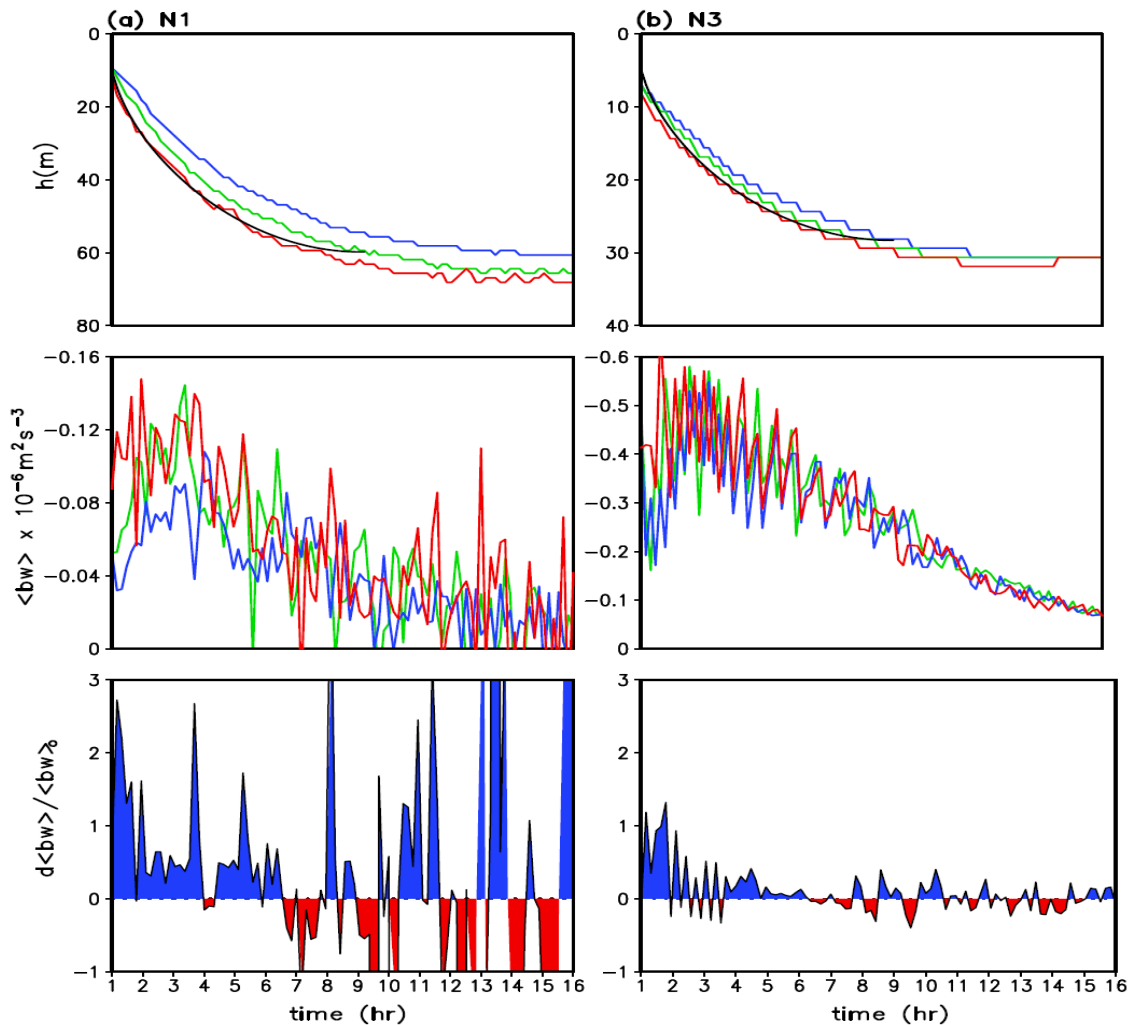
- $u_* = 0.02$ m/s

- $h_0 = 5$ m

- $f = 10^{-4}$ s⁻¹

-LC: a (m) = 0 (L0), 0.5 (L1), 1 (L2)

-Stratification: N^2 (s⁻²) = 10^{-5} (N1), 5×10^{-5} (N2), 2×10^{-4} (N3)



The contribution of LC to the mixed layer deepening is important, only if stratification is weak and MLD is shallow!

Fig. 1 Time series of MLD (h) (upper), the buoyancy flux at the MLD ($\langle bw \rangle$) (middle), and $d\langle bw \rangle / \langle bw \rangle$ (lower). Here $\langle bw \rangle$ is the buoyancy flux in the absence of LC ($\langle bw \rangle$). (blue: L0 ($La = \infty$), green: L1 ($La = 0.64$), red: L2 ($La = 0.32$), black: theoretical prediction by Pollard et al. (1973)): (a) N1 ($\nu = 10^{-5} \text{ s}^{-2}$), (b) N3 ($\nu = 2 \times 10^{-4} \text{ s}^{-2}$).

Dimensional Analysis

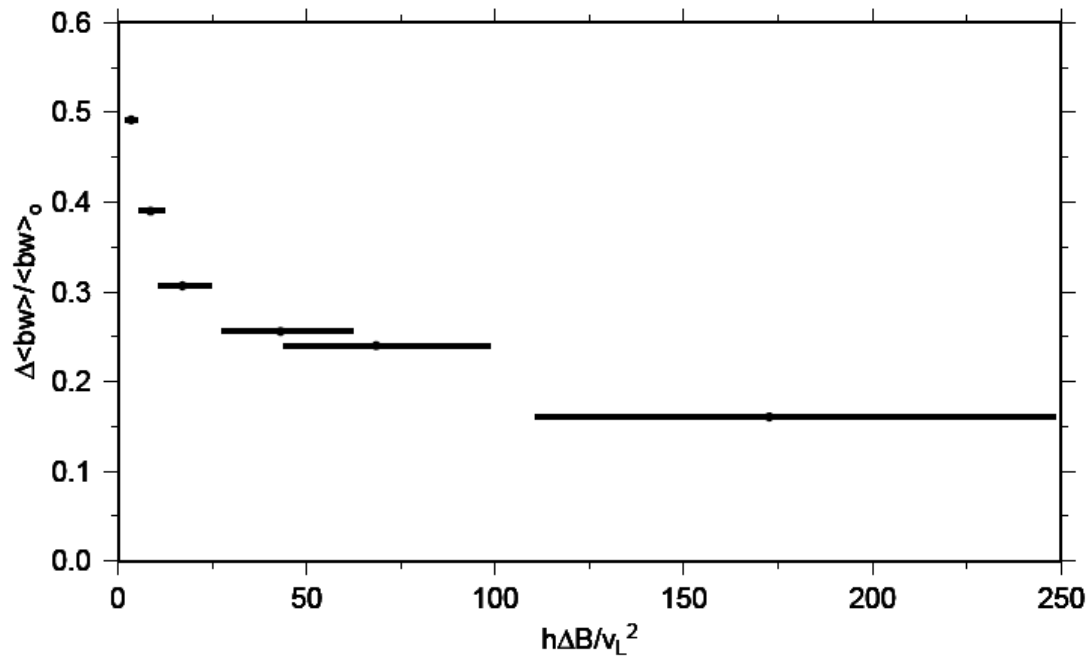
$$w_e = f(h, \Delta B, u_*, U_s, \lambda, f)$$

$$\frac{w_e}{u_*} = f\left(\frac{h\Delta B}{u_*^2}, \frac{h\Delta B}{v_L^2}, \frac{h}{\lambda}, \frac{h}{u_* / f}\right) \quad \leftarrow \quad \begin{aligned} v_L &= (U_s u_*^2)^{1/3} \\ h\Delta B / v_L^2 &= (h\Delta B / u_*^2)(u_* / v_L)^2 \end{aligned}$$

i) If h/λ and $h/(u_*/f)$ remain invariant,

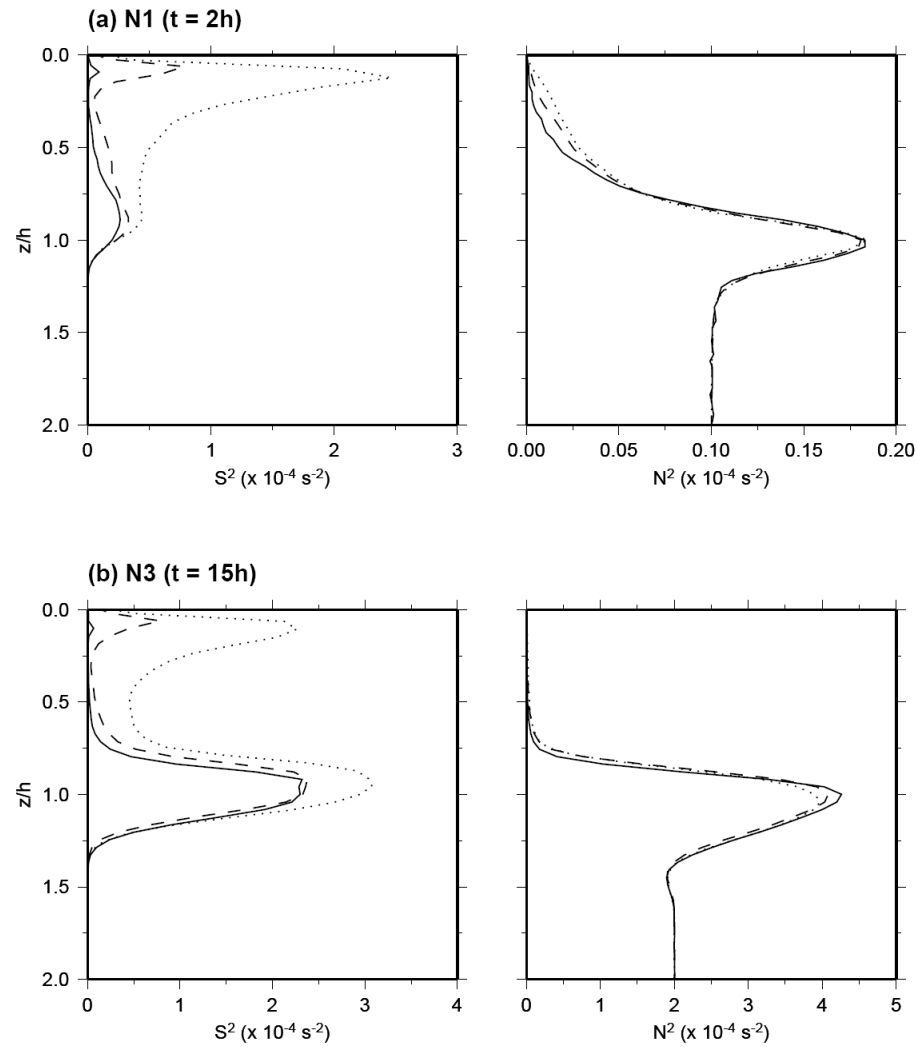
ii) If $w_e / u_* = f_1(h\Delta B / u_*^2)[1 + f_2(h\Delta B / v_L^2)]$

$$\frac{\overline{\Delta bw}(h)}{\overline{bw}(h)_0} = f_2\left(\frac{h\Delta B}{v_L^2}\right) \quad \leftarrow \quad \overline{\Delta bw}(h) = \Delta w_e \Delta B$$



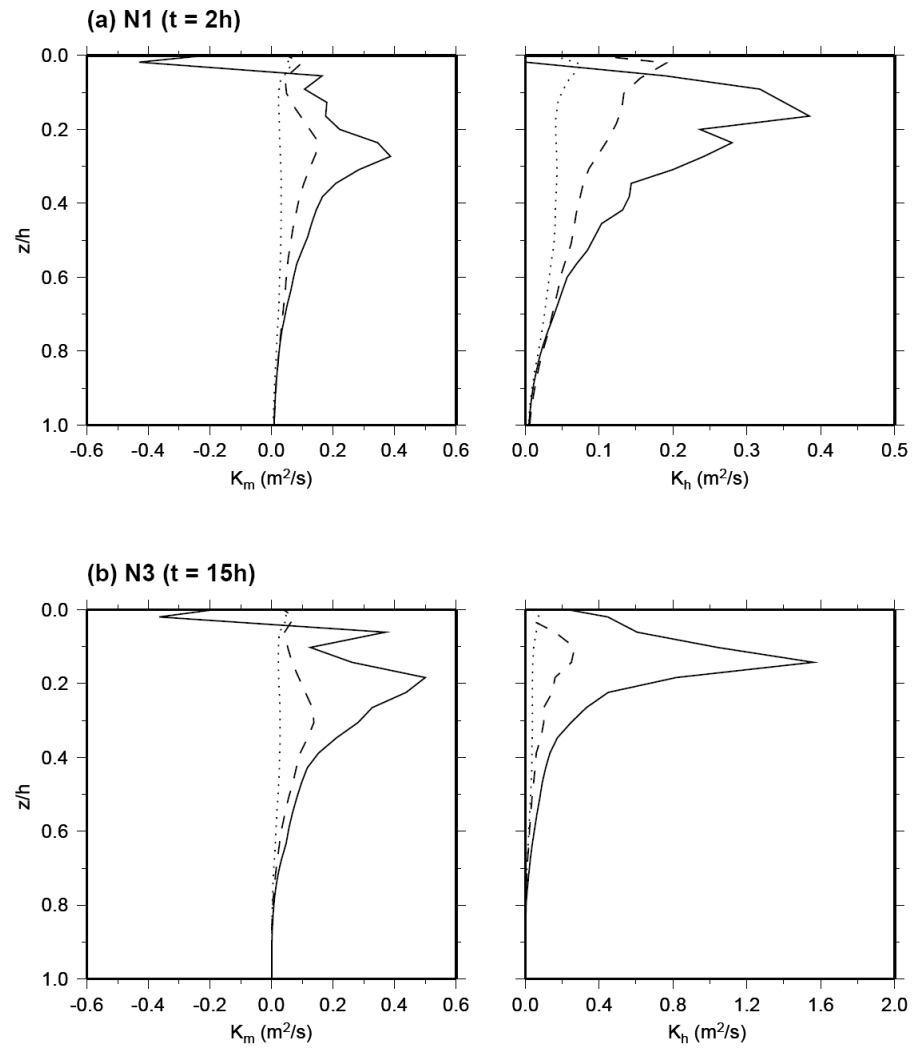
Variation of $\overline{\Delta bw}(h) / \overline{bw}_0(h)$ with $h\Delta B / v_L^2$

⇒ LC induces a significant enhancement of the ML deepening, only if stratification is weak and MLD is shallow.



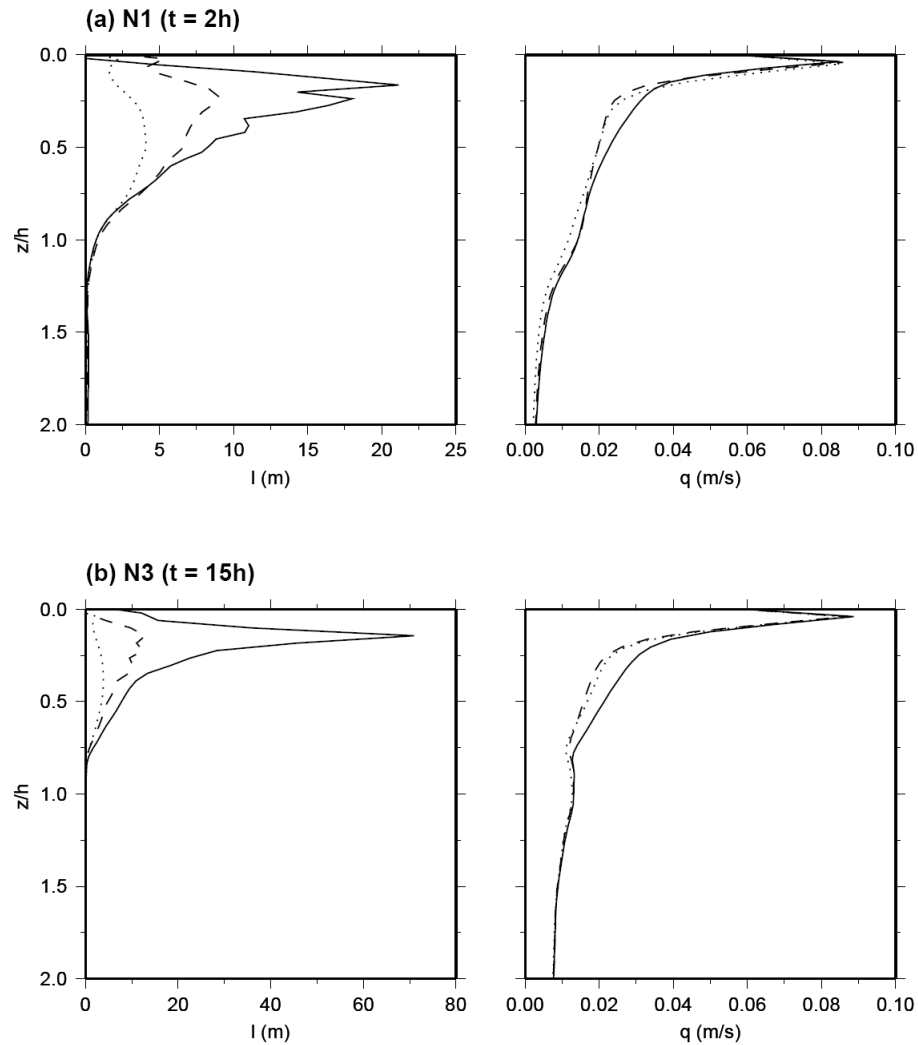
Profiles of S^2 and N^2 (dotted: L0, dashed: L1, solid: L2)

\Rightarrow *much weaker S^2 in the presence of LC*



Profiles of K_m and K_h (dotted: L0, dashed: L1, solid: L2)

⇒ *much larger K_m and K_h within the ML in the presence of LC*

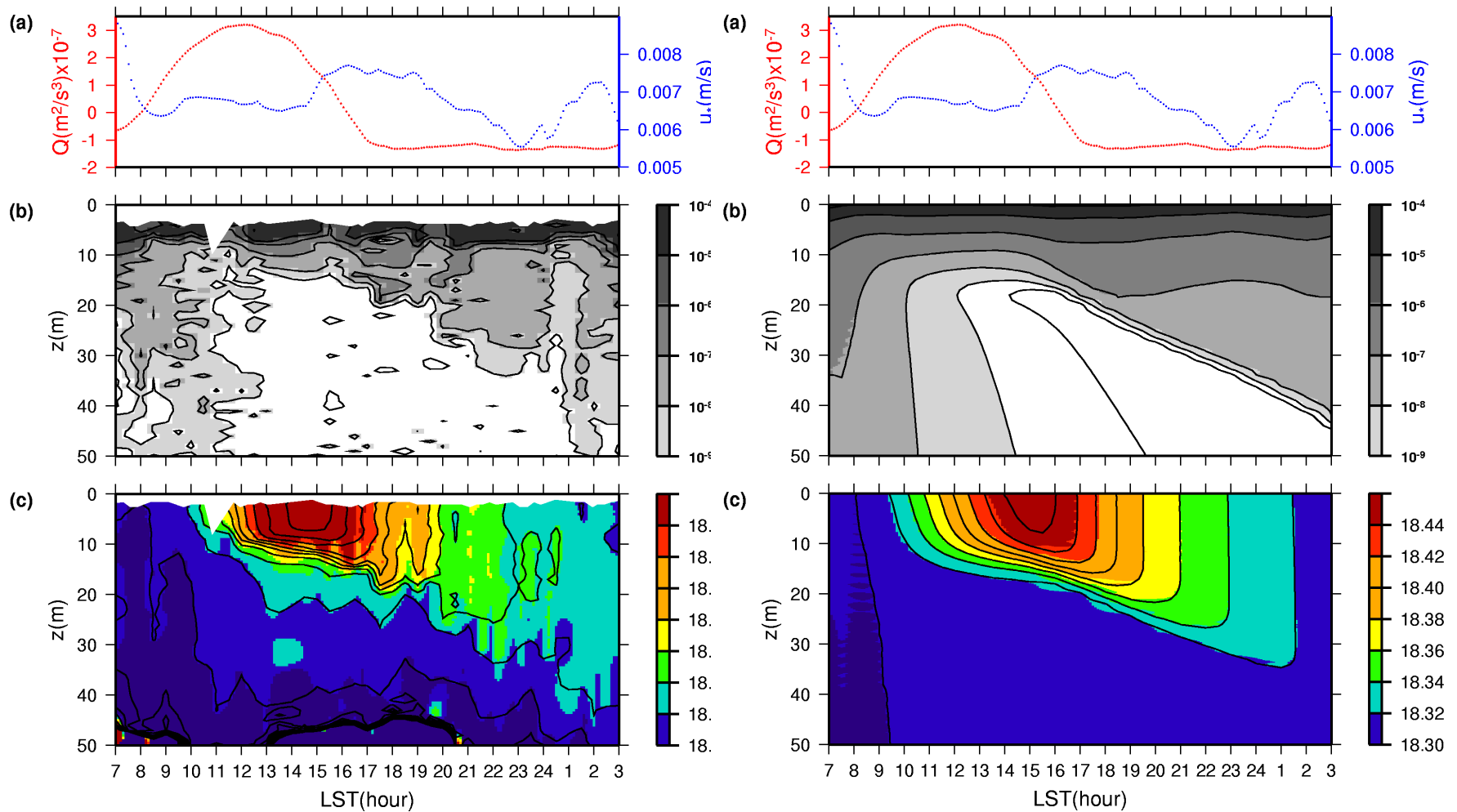


Profiles of l and q (dotted: L0, dashed: L1, solid: L2)

⇒ Much larger K_m and K_h under LC is due to much larger l ($K \sim ql$).

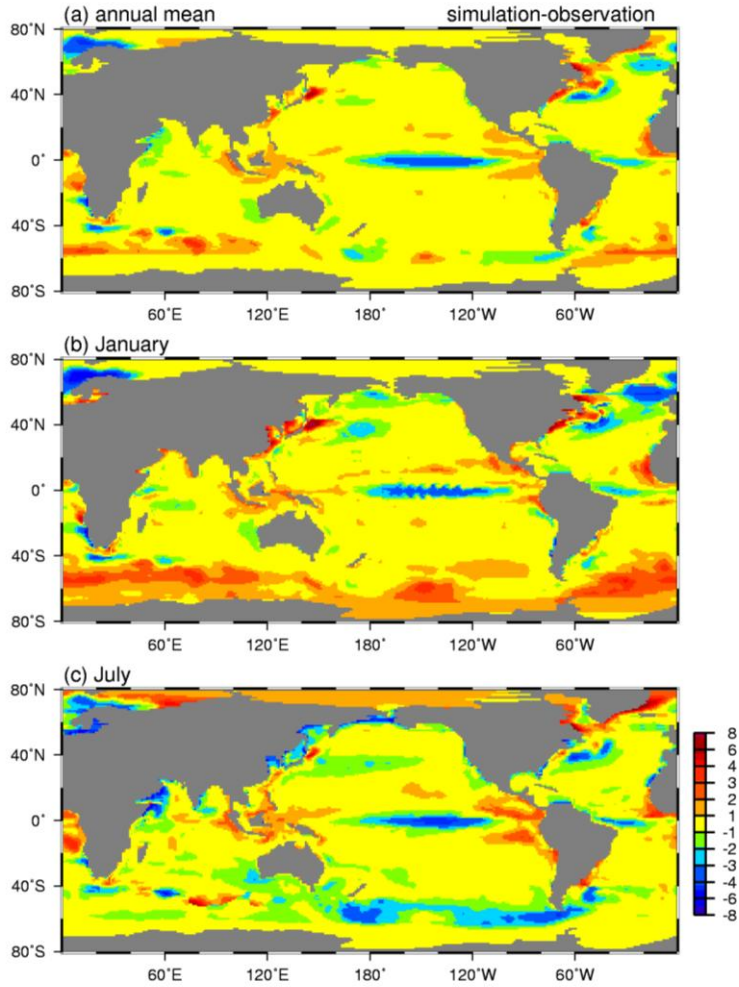
Noh Mixed Layer Model

(Noh et al., *JGR* 1999, *JPO* 2002, *GRL* 2005)

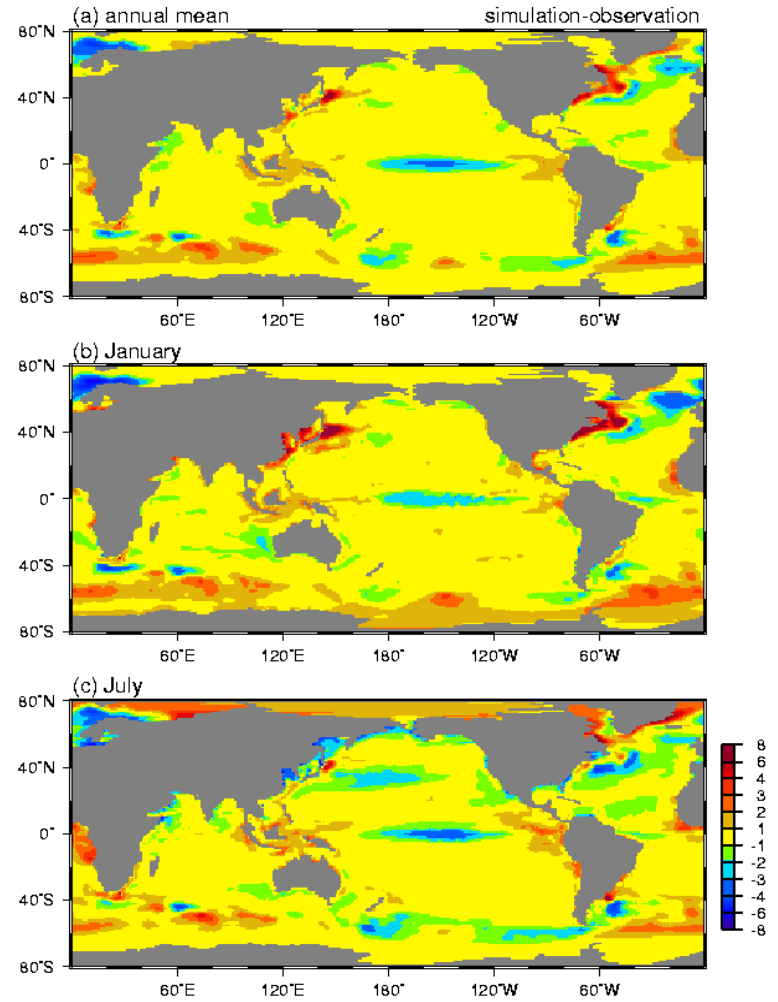


Evolution of the OML from the observation (PATCHEX) (left) and simulation (right): (a) surface forcing, (b) dissipation rate ($\text{m}^2 \text{ s}^{-3}$) and (c) temperature ($^{\circ}\text{C}$)

SST anomaly



PP mixing



a new mixed layer model

Basic Concept of the Noh Model

- **TKE flux at the surface;**

$$K_E \partial E / \partial z = Au_*^3 \quad \text{with } A = 100$$

**Mellor & Yamada (1982): $A = 0$ m - neglect wave breaking*

- **effects of stratification;**

$$l/l_0 = (1 + \alpha \text{Rt})^{-1/2} \quad \leftarrow \quad K \sim ql$$

$$\text{Rt} = (Nl_0 / q)^2$$

**Mellor & Yamada (1982), Pacanowski & Philander (1981)*

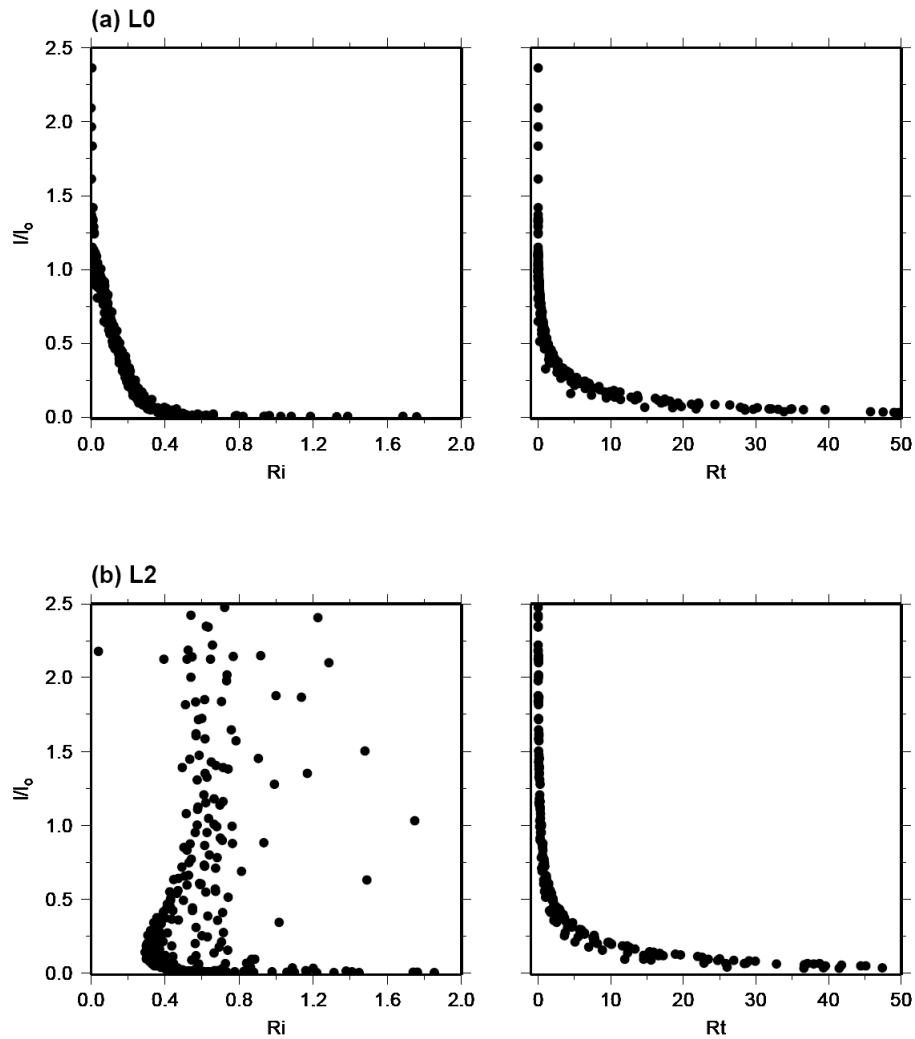
$$l/l_0 = f(\text{Ri}) \quad \text{Ri} = (N/S)^2$$

- TKE production is dominated by shear production ($P_s - P_b - \varepsilon = 0$).

- **length scale;**

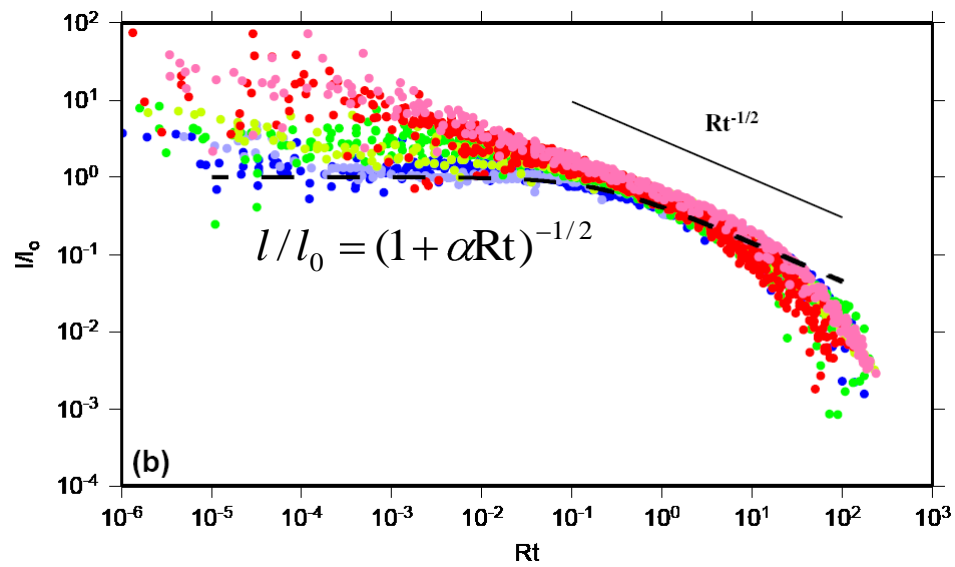
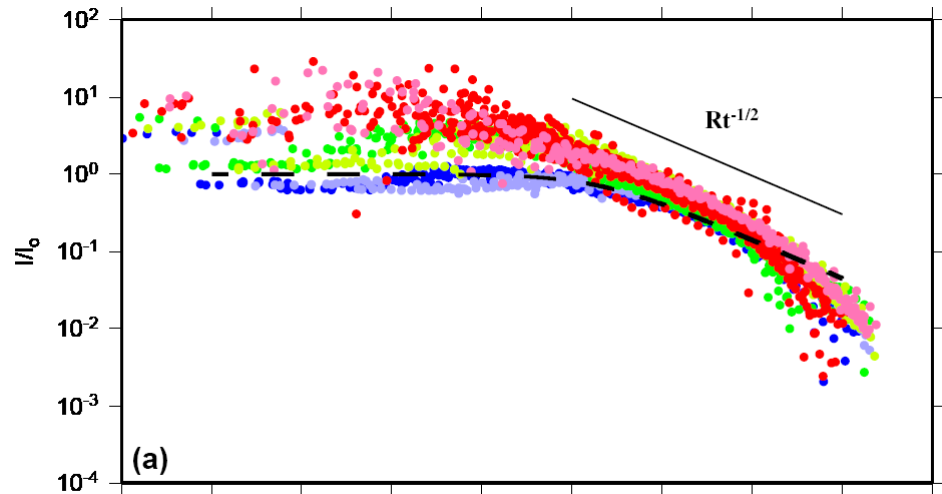
$$\frac{1}{l_0} = \frac{1}{\kappa(z + z_0)} + \frac{1}{h} \quad \text{with } z_0 = 1 \text{ m}$$

**Mellor & Yamada (1982): $z_0 = 0$ m - neglect wave breaking*



Scatter plots l/l_0 vs. Ri and Rt (N3): (a) L0, (b) L2

$\Rightarrow Ri$ cannot be a parameter in the presence of LC



L0: blue, L1: green, L2: red

Conclusion

- LC induces a significant enhancement of the mixed layer deepening, only if stratification is weak and MLD is shallow.
- Stronger vertical mixing within the ML in the presence of LC is due to much larger l .
- Vertical mixing must be parameterized in terms of Rt instead of Ri in the presence of LC.
- Vertical mixing in the presence of LC can be parameterized by

$$\Rightarrow \quad l/l_0 = \gamma(1 + \alpha\gamma^2 Rt)^{-1/2} \quad (\alpha \sim 5, \gamma \sim 10)$$

2. What is the Role of **Wind Stress** in the Convective Deepening of the Mixed Layer ?

Parameterization for convective deepening

© NK model [Niiler and Kraus, 1977]

$$-\overline{bw}_h = -nQ_0 + 2m_0 u_*^3 / h + m_s w_e \overline{U^2} / h \quad (n = 0.21, m_0 = 0.39, m_s = 0.48)$$

$$* \quad -\overline{bw}_h = w_e \Delta B, \quad w_e = \partial h / \partial t$$

© PWP model [Price et al., 1986]

$$Ri_b = \frac{[B_0 - B(h)]h}{|\mathbf{U}_0 - \mathbf{U}(h)|^2}, \quad Ri_g = \frac{\partial B / \partial z}{(\partial \mathbf{U} / \partial z)^2} = \text{const. at the MLD}$$

© KPP model [Large et al., 1994]

$$Ri_b^* = \frac{[B_0 - B(h)]h}{|\mathbf{U}_0 - \mathbf{U}(h)|^2 + V_t^2} = \text{const. at the MLD}$$

* Q_0 = surface buoyancy flux, ρu_*^2 = wind stress, V_t = convective velocity scale

Simulations

@ model domain

@ $L_x = L_y = 300 \text{ m}$ & $H = 80 \text{ m}$

@ $\Delta x = \Delta y = \Delta z = 1.25 \text{ m}$ with $n_x = n_y = 240$ & $n_z = 64$

@ forcing

@ $Q_0 = 1.25, 2.5, 5 \times 10^{-7} \text{ m}^2/\text{s}^3$

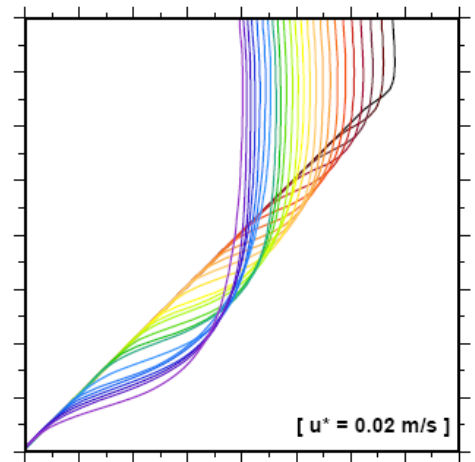
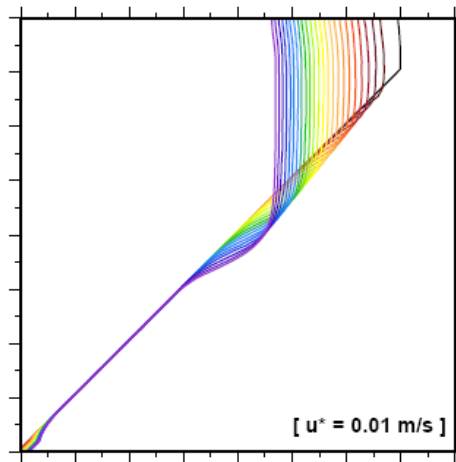
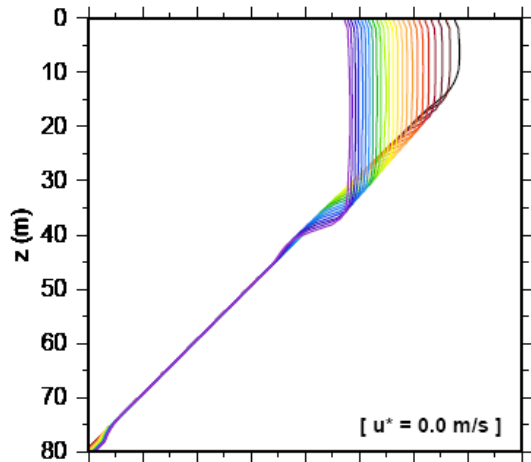
@ for rotation effect : $\varphi = 0^\circ, 30^\circ, 55^\circ$

@ for wind stress : $u^* = 0.0, 0.01, \& 0.02 \text{ m/s}$

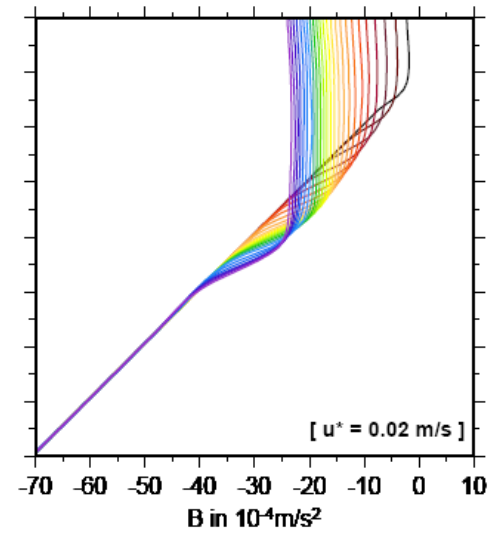
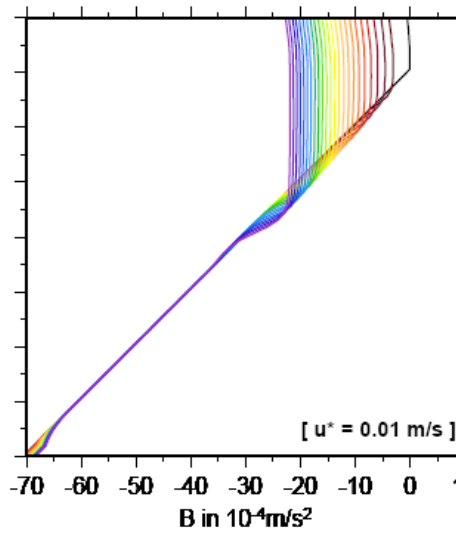
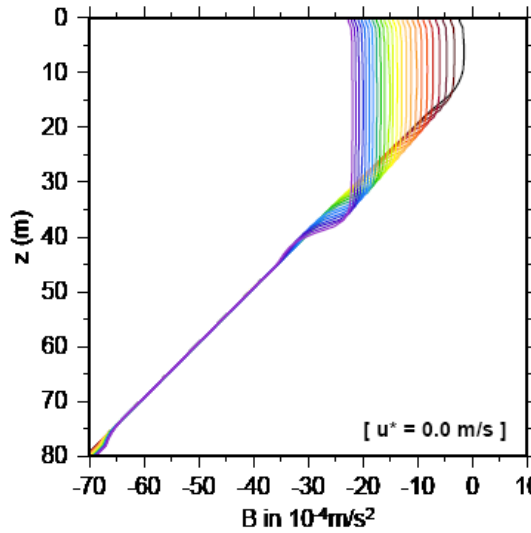
@ 12 hours integration

Buoyancy

$\varphi = 0^\circ$



$\varphi = 55^\circ$

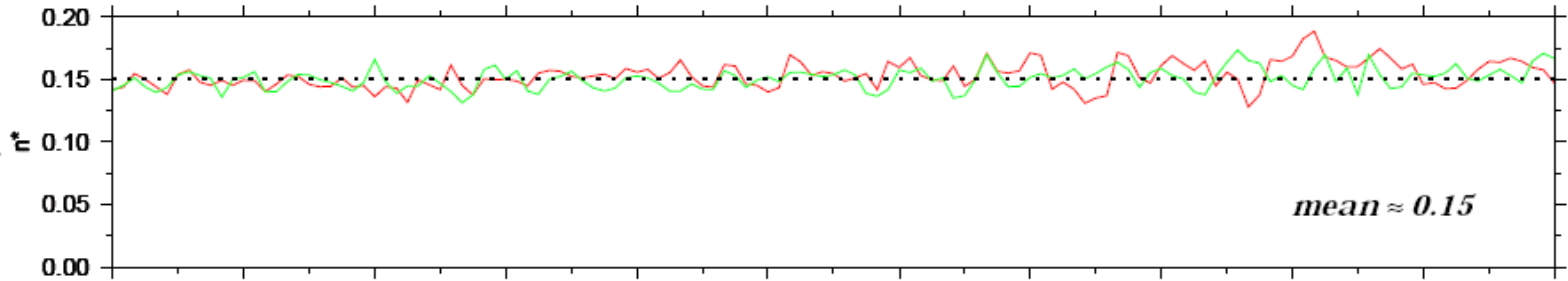


$$-\overline{bw}_h = -nQ_0 + 2m_0u_*^3/h + m_s w_e \Delta U^2/h \quad (\text{NK model})$$

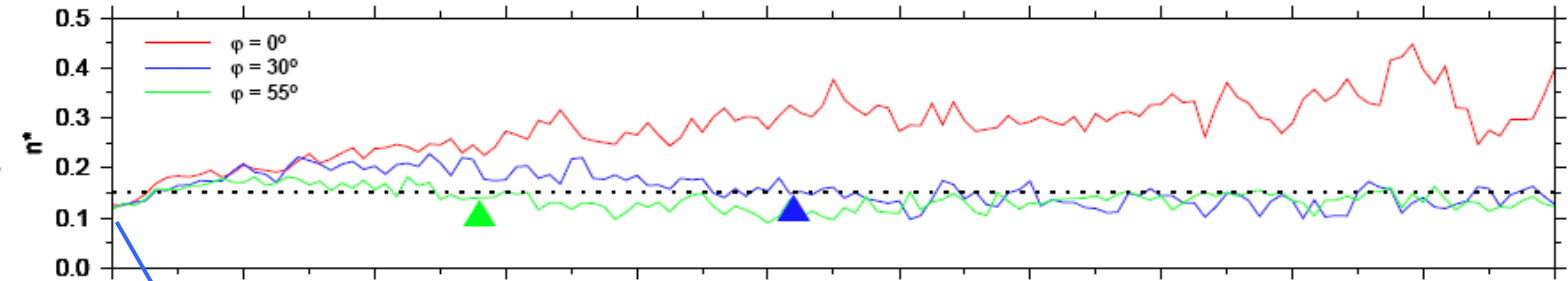
$$-\overline{bw}_h = -n^*w_*^3/h, \quad w_*^3 = Q_0h$$

▲ : $t \sim \pi/f$

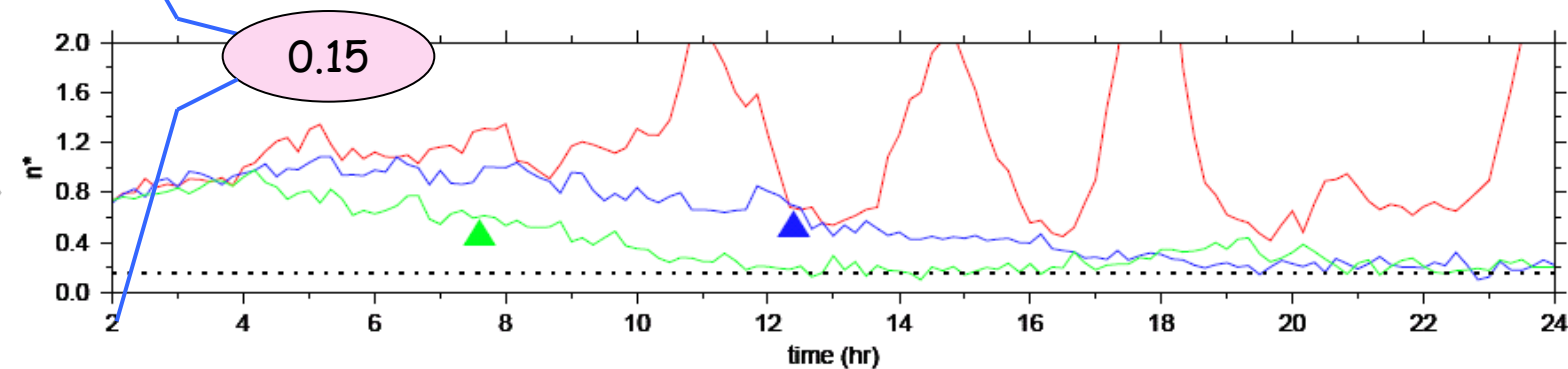
$u^* = 0.0$
m/s



$u^* = 0.01$
m/s

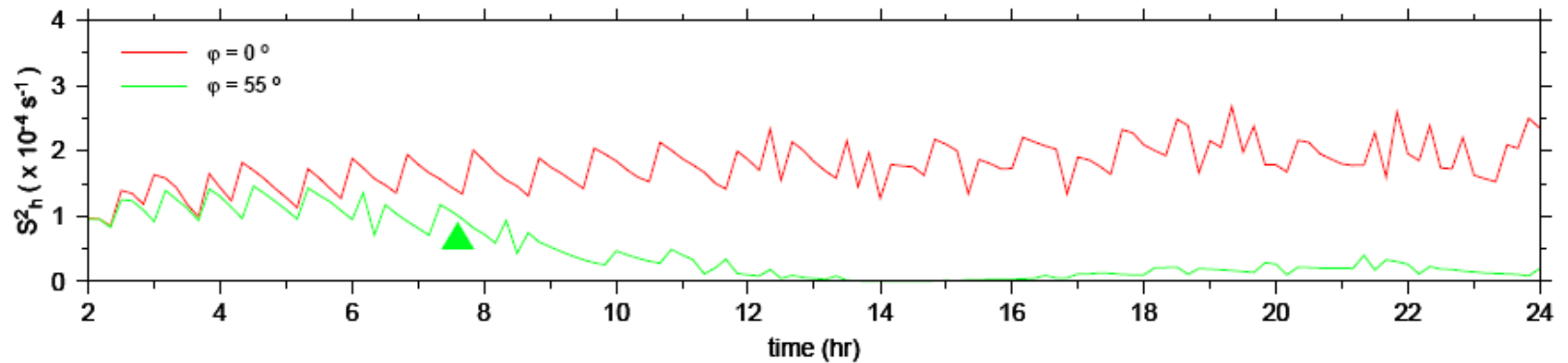
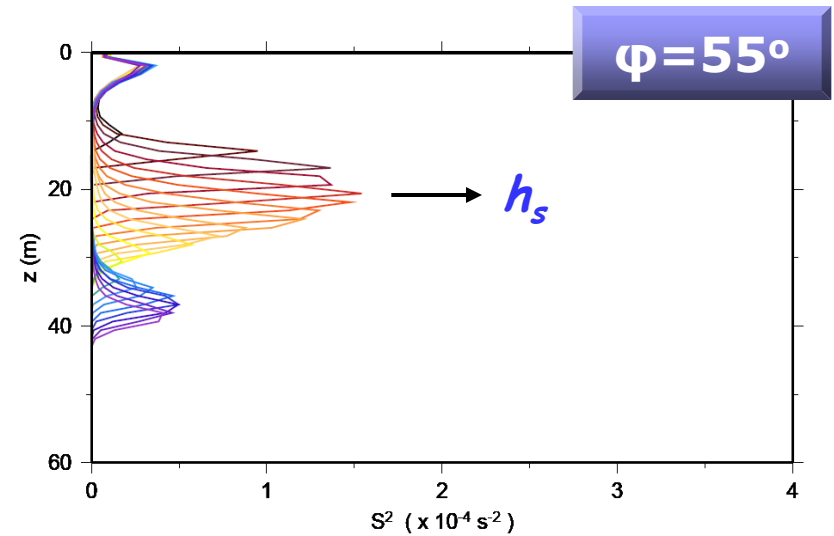
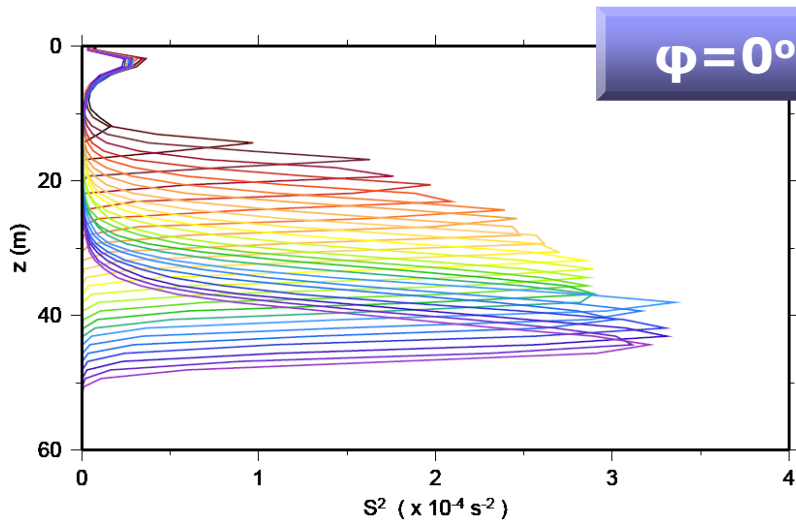


$u^* = 0.02$
m/s

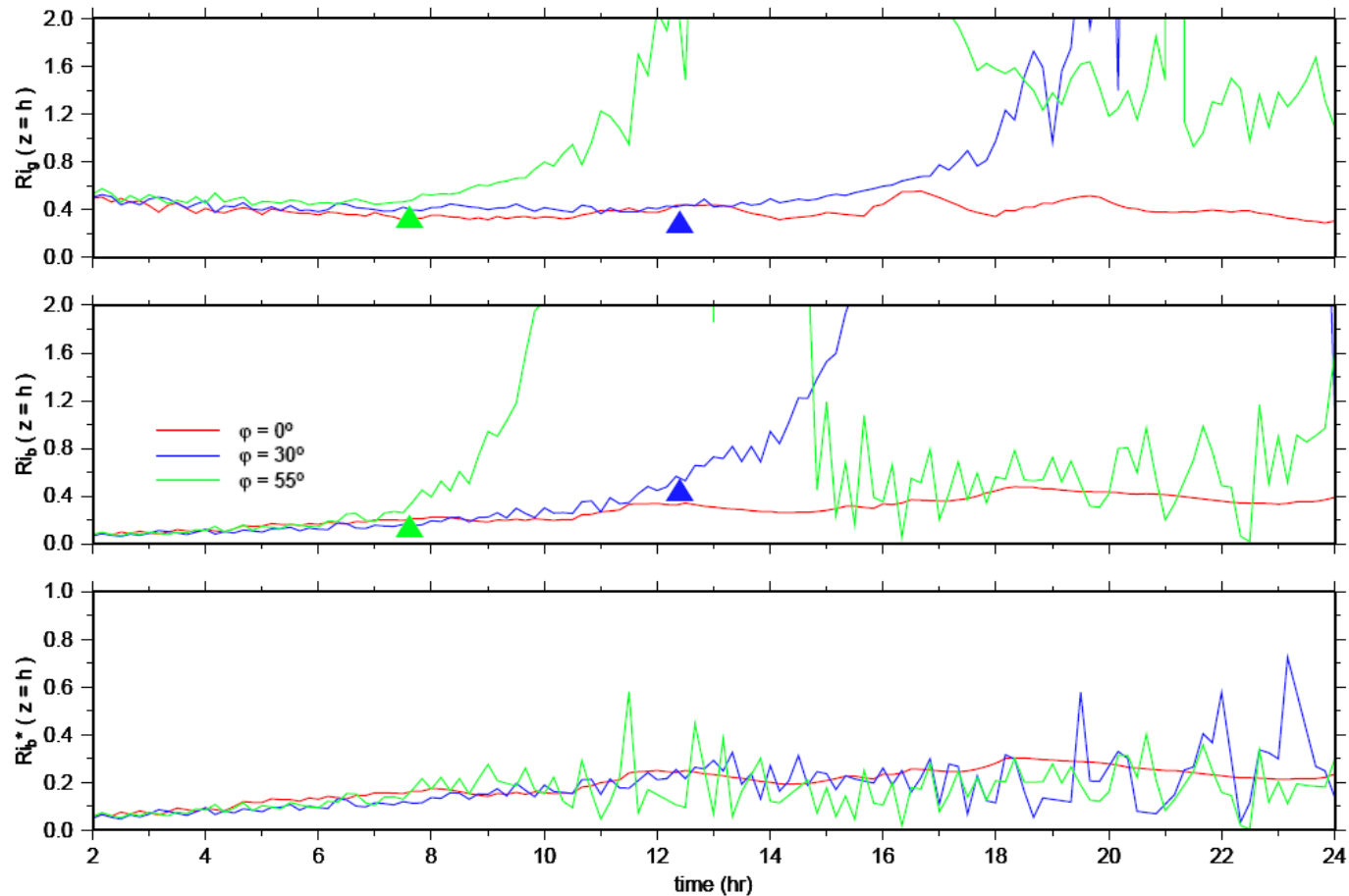


➔ *The effects of wind forcing m_0 in the NK model is overestimated in the high latitude!*

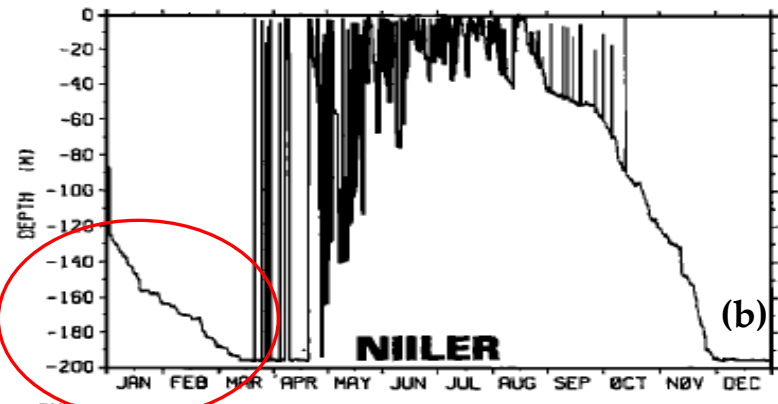
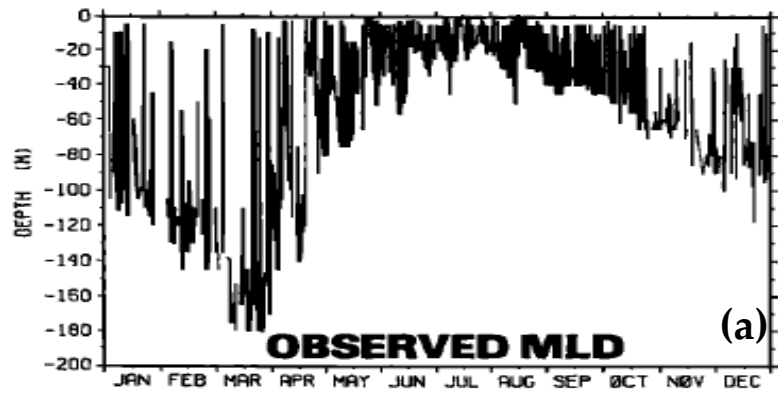
$$S^2 = \left(\frac{\partial u}{\partial z} \right)^2 + \left(\frac{\partial v}{\partial z} \right)^2, \quad h_s = \text{the depth of maximum shear}$$



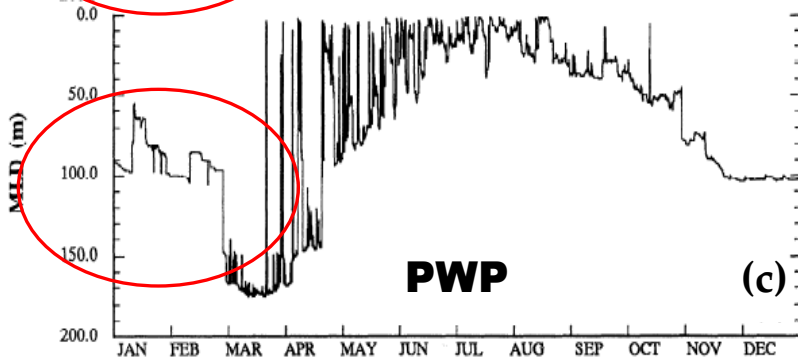
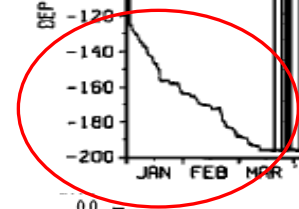
⊙ Ri_g , Ri_b , and Ri_b^* ($z = h$) for $u^* = 0.02$ m/s



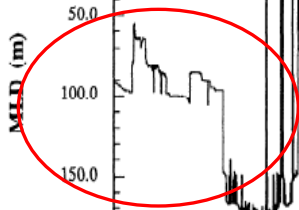
→ *The critical Ri_g and Ri_b are overestimated in the PWP model!
The critical Ri_b^* in the KPP model is OK!*



overestimation



underestimation

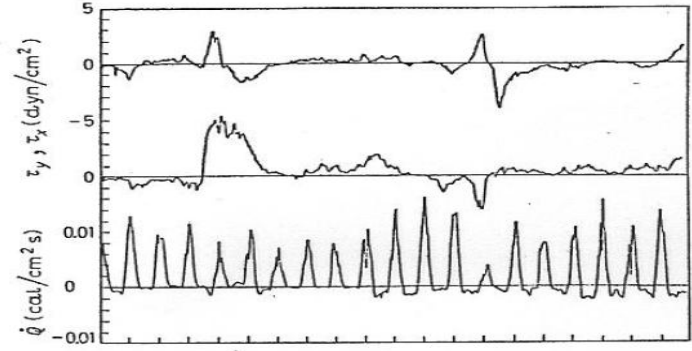
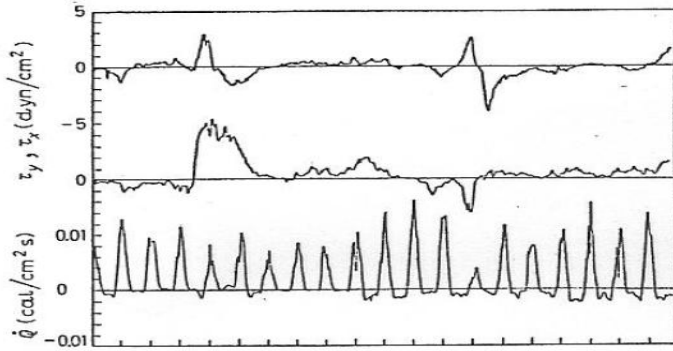


Seasonal variation of MLD simulated by NK and PWP model (Station PAPA 1961; 50 N)

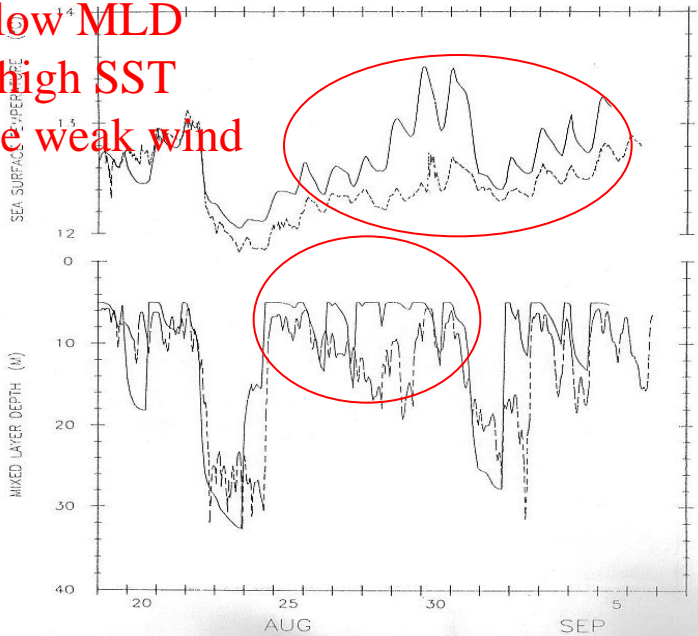
Sensitivity Test of NK model

@ NK model

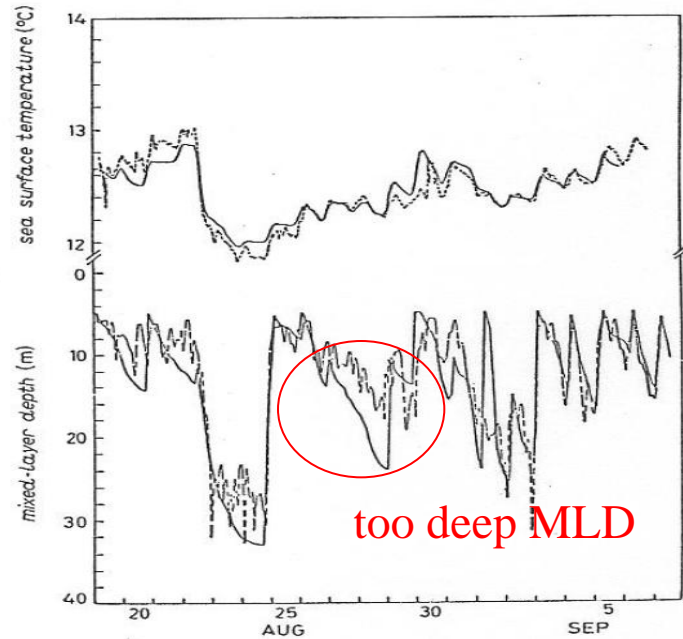
[Davis et al., 1981]



too shallow MLD
and too high SST
under the weak wind



($m_0 = 0.13$, $m_s = 0.67$)



too deep MLD

($m_0 = 0.39$, $m_s = 0.48$)

Suggested Modification of NK model

$$-\overline{bw}_h = nQ_0 + \boxed{2m_0 u_*^3 e^{z/\delta} / h} + m_s w_e (\Delta U)^2 / h$$

effect of wave breaking and Lanmuir circulation in the near surface region

Conclusion

- In the extratropical ocean,
 - the contribution from wind stress disappears after the inertial period
 - the bulk and gradient Richardson number at the MLD increases rapidly with time.

→ NK model and PWP models cannot be applied.
- In the tropical ocean,
 - the contribution from wind stress continues to increase until it reaches a critical value.

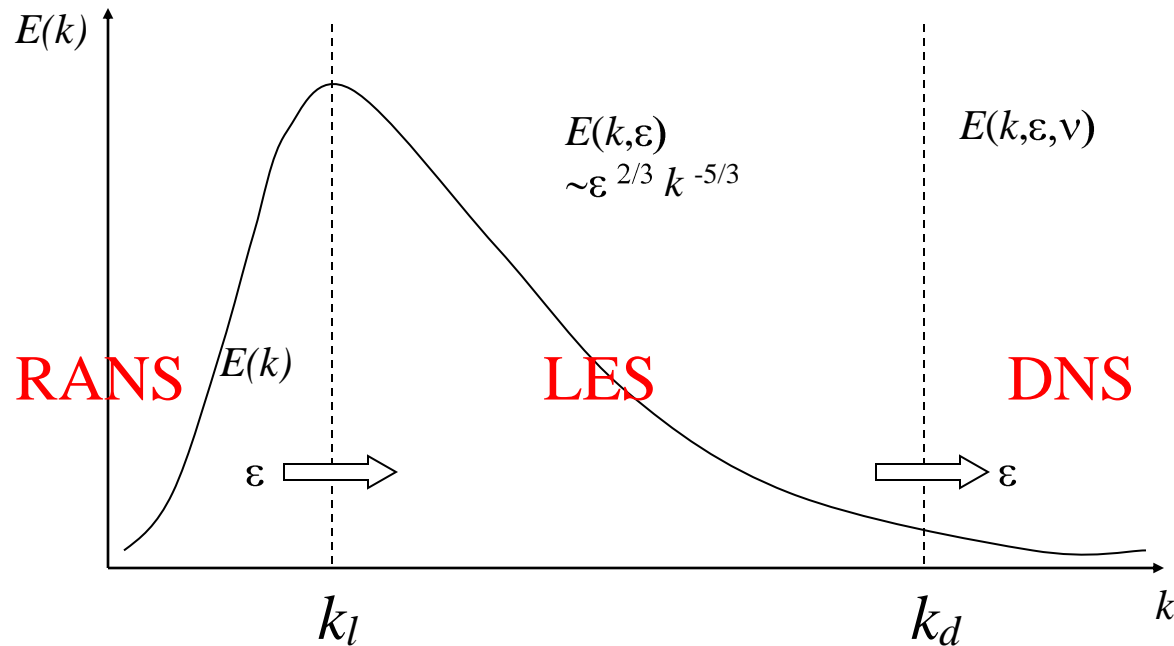
Numerical Ocean Circulation Modelling: Present Status and Future Direction (J. Willebrand & D.B. Haidvogel, in *Ocean Circulation and Climate*)

*"The most pressing concern of the OGCM, particularly for global climate modeling, is the improvement of **parameterization for SGS process**"*

*"Especially, **the surface mixed layer** has received special attention because of its important role in air-sea exchange"*

*"A strong test of a SGS closure would involve a high-resolution simulation such as **LES**, followed by a demonstration that a coarse resolution OGCM with the new parameterization gives qualitatively identical net results."*

What is LES?



- reproduces the realistic 3-D structure of fluctuating turbulent flows by utilizing strong computing power

- Most turbulence models of the OML apply the same assumptions and parameterizations used for the ABL.

- Mellor and Yamada (1982), Kantha and Clayson (1994), Large et al. (1994), Klein and Coantic (1981), Andre and Lacarrere (1985), Mellor and Durbin (1975), Canuto (2001), Umlauf et al. (2004)

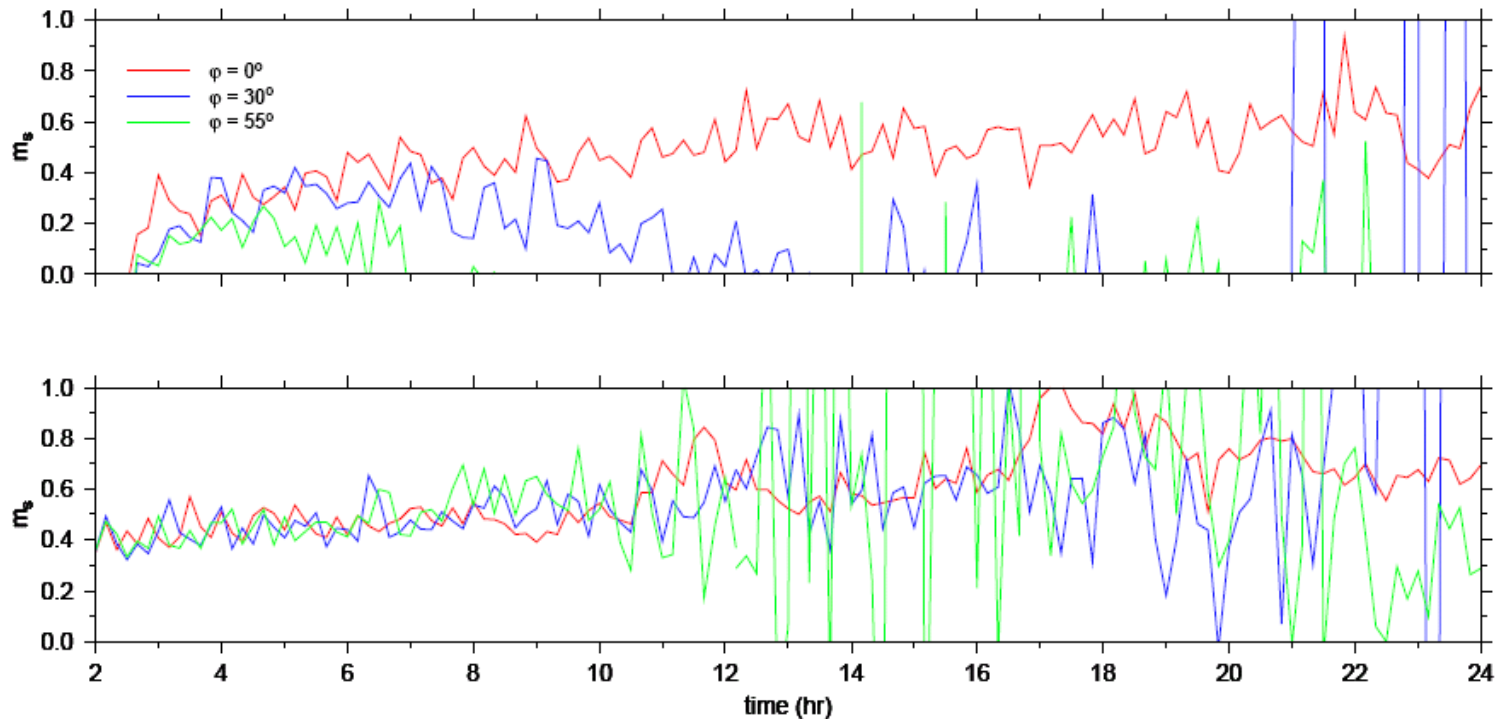
*"In view of the fact that they do not explicitly incorporate the unique process occurring in the oceanic mixed layer, **such as Langmuir circulation and wave breaking**, higher-order closure models of the oceanic mixed layer are **similar to geocentric pre-Copernican astronomy**, in which epicycles were added to epicycles to account for the apparent motions of the planetes" (E. Kunze, C. Garrett)*

$$-\overline{bw}_h = -nQ_0 + 2m_0u_*^3/h + m_s w_e \Delta U^2/h$$

e.g., $n = 0.21$, $m_0 = 0.39$, and $m_s = 0.48$ [Davis et al., 1981]

$m_0 = 0$ with $n = 0.15$

$u_* =$
0.01 m/s



$u_* =$
0.02 m/s